

# Crystal structure of $(\text{Ba}_{1-x}\text{La}_x)$ $[\text{Mg}_{(1+x)/3}\text{Nb}_{(2-x)/3}]\text{O}_3$ ceramics

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The crystal structure of  $(\text{Ba}_{1-x}\text{La}_x)[\text{Mg}_{(1+x)/3}\text{Nb}_{(2-x)/3}]\text{O}_3$  (BLMN) ceramics with  $0 \leq x \leq 1$  was investigated using X-ray diffraction (XRD) and high resolution transmission electron microscopy (HRTEM). When the La content,  $x$ , was above 0.1, the 1:2 ordered hexagonal structure found in  $\text{Ba}(\text{Mg}_{1/3}\text{Nb}_{2/3})\text{O}_3$  (BMN) was transformed into the 1:1 ordered cubic structure. The 1:1 ordered cubic structure was maintained up to  $x = 0.7$ . When  $x$  exceeded 0.7, however, BLMN exhibited a 1:1 ordered monoclinic structure, rather than a 1:1 ordered cubic structure.  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$  (LMN) has a 1:1 ordered monoclinic  $P2_1/n$  structure with  $a = 5.6004 \text{ \AA}$ ,  $b = 5.6414 \text{ \AA}$ ,  $c = 7.9346 \text{ \AA}$ , and  $\beta = 89.9819^\circ$ . The monoclinic LMN has the in-phase and the anti-phase tilt of oxygen octahedra. The anti-parallel shift of A-site cations was also found in LMN. © 2003 Kluwer Academic Publishers

## 1. Introduction

Complex perovskite ceramics have been extensively investigated for the application to the dielectric resonators at microwave frequencies [1]. The dielectric properties of these materials are largely affected by the structural properties. In particular, the chemical ordering of B site cations and the tilting of the oxygen octahedra are the major factors which influence  $Q$  value and the temperature coefficient of resonance frequency ( $\tau_f$ ) of these ceramics. According to the previous investigations, the 1:2 chemical ordering improved the  $Q$  value of the complex perovskite ceramics [2]. Tilting of the oxygen octahedra in simple perovskite was investigated by Glazer [3, 4] and Aleksandrov [5, 6]. Based on the Glazer's model, Colla and Reaney have studied the correlation between  $\tau_f$  and octahedra tilt in the Ba- and Sr-based 1:2 ordered complex perovskite systems [7, 8].

$\text{Ba}(\text{Mg}_{1/3}\text{Nb}_{2/3})\text{O}_3$  ceramics, one of the complex perovskites, have high values of  $\varepsilon_r$  (32), and  $Q$  value ( $Q * f_0 = 56,000$ ) [9]. These ceramics, however, have relatively high  $\tau_f$  (33 ppm/°C) compared with those of other complex perovskites [2, 9, 10]. The temperature coefficient of resonant frequency of BMN can be tailored by the addition of the materials with a negative  $\tau_f$  value.  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$  has the value  $\tau_f = -86 \text{ ppm/}^\circ\text{C}$ ,  $\varepsilon_r = 26.8$  and  $Q * f_0 = 52,000$  [11]. A solid solution of BMN and LMN there-

fore is expected to have zero  $\tau_f$  with high  $Q$ -value and  $\varepsilon_r$ . However, the crystal structures of BLMN and LMN have not been studied. Rhombohedral  $R\bar{3}$  or orthorhombic  $Pbmn$  structure was proposed for LMN [12, 13]. However, according to our analysis, they are not the correct structures for LMN. In this work,  $(\text{Ba}_{1-x}\text{La}_x)[\text{Mg}_{(1+x)/3}\text{Nb}_{(2-x)/3}]\text{O}_3$  ceramics with  $0 \leq x \leq 1$  were prepared, and the structural transition of BLMN with La content was studied using HRTEM and XRD.

## 2. Experimental procedures

Specimens of  $(\text{Ba}_{1-x}\text{La}_x)[\text{Mg}_{(1+x)/3}\text{Nb}_{(2-x)/3}]\text{O}_3$ , with  $0 \leq x \leq 1$  were synthesized by a conventional solid-state reaction method from  $\text{BaCO}_3$  (Aldrich, 99%),  $\text{La}_2\text{O}_3$  (Aldrich, 99%),  $\text{Nb}_2\text{O}_5$  (tritech, 99.9%), and  $\text{MgO}$  (Junsei, 99%). The materials were weighed in the appropriate molar ratio and mixed with  $\text{ZrO}_2$  balls in an ethanol media for 24 h. The powders were dried, calcined at  $1300^\circ\text{C}$  for 4 h in air, and pressed into a disk. These pellets were sintered at  $1450^\circ\text{C}$  for 4 h in air. X-ray diffraction pattern (XRD) data for the Rietveld refinement were collected at room temperature with MACScience M18XHF diffractometer. Rietveld refinements were performed using DBWS 9411 [14]. Microstructure of the specimens were studied using high resolution transmission electron

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microscopy (HRTEM). TEM specimens were cut from disc, mechanically ground, and ion milled using argon. Hitachi H-9000NAR TEM were used to examine the specimens.

### 3. Results and discussion

Fig. 1 shows the X-ray powder diffraction patterns of  $(\text{Ba}_{1-x}\text{La}_x)[\text{Mg}_{(1+x)/3}\text{Nb}_{(2-x)/3}]\text{O}_3$ , with  $0 \leq x \leq 1$  specimens sintered at  $1450^\circ\text{C}$  for 4 h. All the peaks were indexed based on the cubic perovskite unit cell. As seen from the figure, the intensities of the 1:2 ordering superlattice reflections shown in BMN decreased as La content,  $x$ , increased, eventually disappearing. When  $x > 0.1$ , a series of extra peaks at  $1/2(111)$  and  $1/2(311)$  positions — indicated by open circle — began to grow, attaining maxima at  $x = 0.5$ . When  $x$  was above 0.5, the intensity of  $1/2(111)$  extra peak gradually decreased but that of  $1/2(311)$  extra peak has not been changed. The extra peak at  $1/2(111)$  position has formed as a result of the 1:1 chemical ordering of cations. Because two types of ions existed in both the A and B sites, the 1:1 chemical ordering could occur in the A sites and/or B sites. However, since the 1:1 ordering also was found in LMN which had only La ions in the A site, the 1:1 ordering was considered restricted in B sites. The same results were found by other investigators for the systems  $(\text{Pb}_{1-x}\text{La}_x)[\text{Mg}_{(1+x)/3}\text{Nb}_{(2-x)/3}]\text{O}_3$  and  $(\text{Ba}_{1-x}\text{La}_x)[\text{Zn}_{(1+x)/3}\text{Ta}_{(2-x)/3}]\text{O}_3$  [15, 16].

The  $1/2(311)$  extra peak shown in Fig. 1 could arise from the 1:1 chemical ordering and/or the anti-phase tilt of octahedra [3, 4, 17]. When  $x < 0.7$ , the variation of the intensity of  $1/2(311)$  with  $x$  was the same as that of  $1/2(111)$  extra peak thus, the  $1/2(311)$  extra peak was thought to be formed due to the 1:1 ordering. On the contrary, when  $x > 0.7$ , the degree of the 1:1 ordering identified by the  $1/2(111)$  extra peak decreased but the intensity of  $1/2(311)$  extra peak has not been changed. This result implies that when  $x$  exceeded 0.7, the  $1/2(311)$  extra peak was associated with the anti-phase tilting of the octahedra as well as the 1:1 chemi-

cal ordering. More discussion on the anti-phase tilting of the octahedra will be given in later.

As seen in Fig. 1, BLMN had a 1:1 ordered cubic structure when  $0.1 < x < 0.7$ . As  $x$  exceeded 0.7, however, new extra peaks appeared at the  $1/2(210)$ ,  $1/2(300)$  and  $1/2(410)$  positions, indicated by closed circles. The anti-parallel shift of the A-site cations is responsible for the presence of those extra peaks [8]. Therefore, when  $x > 0.7$ , the 1:1 ordered cubic is not the structure of BLMN.

In order to determine the crystal structure of BLMN with  $x > 0.7$ , the detailed investigation on the crystal structure of LMN has been conducted. Fig. 2 shows the XRD pattern of  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$  ceramics sintered at  $1450^\circ\text{C}$  for 4 h. As can be seen in this figure, a series of extra peaks appeared at the  $1/2(111)$ ,  $1/2(210)$ ,  $1/2(300)$ ,  $1/2(311)$  and  $1/2(410)$  positions. The appearance of the  $1/2(111)$  extra peak was result of the 1:1 chemical ordering of B site ions. The presence of the  $1/2(210)$ ,  $1/2(300)$  and  $1/2(410)$  extra peaks indicates that the anti-parallel shift of the A-site cations exists in LMN. The  $1/2(311)$  extra peak could arise from the 1:1 chemical ordering and/or the anti-phase tilt of octahedra. In general, the anti-phase tilting of the octahedra precedes the anti-parallel shift of the A-site cations [7]. Therefore, it is thought that the anti-phase tilting of the octahedra exists in LMN and both the 1:1 chemical ordering and the anti-phase tilting were responsible for the formation of the  $1/2(311)$  extra peak.

In order to prove the presence of anti-phase tilting of the octahedra in LMN, the solid solution of  $(1-x)\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3-x\text{SrTiO}_3$  was formed. Since the  $\text{SrTiO}_3$  has the cubic structure and only one ion,  $\text{Ti}^{4+}$ , exists in B site, the lattice distortion and the 1:1 chemical ordering existed in LMN were expected to decrease with the increase of  $x$ . Fig. 3 shows the XRD patterns of  $(1-x)\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3-x\text{SrTiO}_3$  system sintered at  $1500^\circ\text{C}$  for 4 h. As can be seen in this figure, the 1:1 ordering identified by the  $1/2(111)$  extra peak decreased with increasing  $x$  and finally disappeared when  $x > 0.3$ . However, the  $1/2(311)$  extra

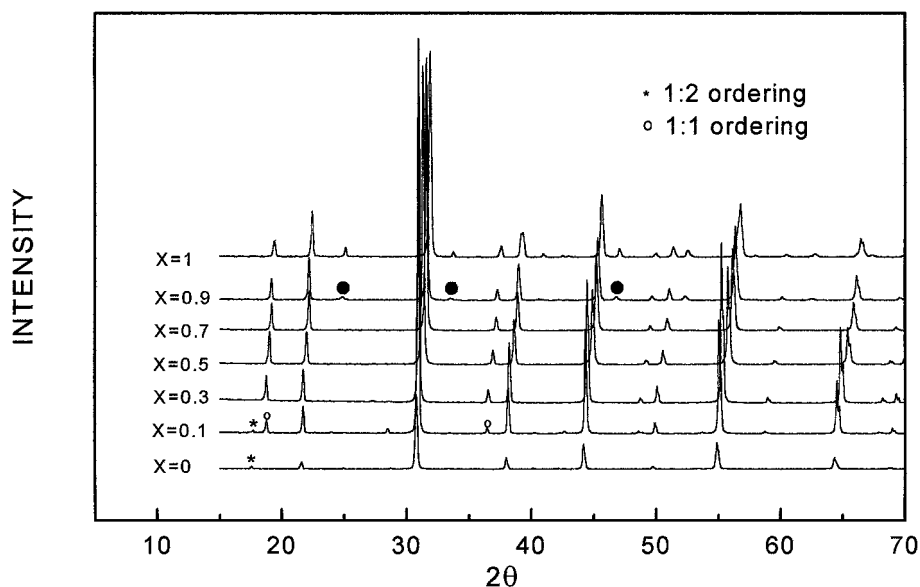


Figure 1 X-ray powder diffraction pattern of  $(\text{Ba}_{1-x}\text{La}_x)(\text{Mg}_{(1+x)/3}\text{Nb}_{(2-x)/3})\text{O}_3$  with  $0 \leq x \leq 1$  sintered at  $1450^\circ\text{C}$  for 4 h.

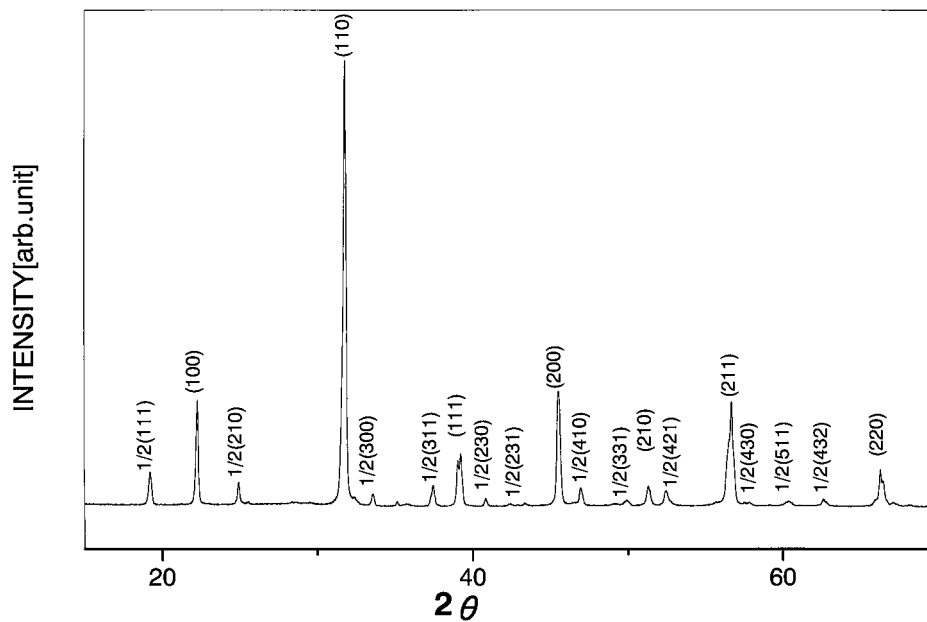


Figure 2 X-ray diffraction pattern of the  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$  ceramics sintered at  $1450^\circ\text{C}$  for 4 h.

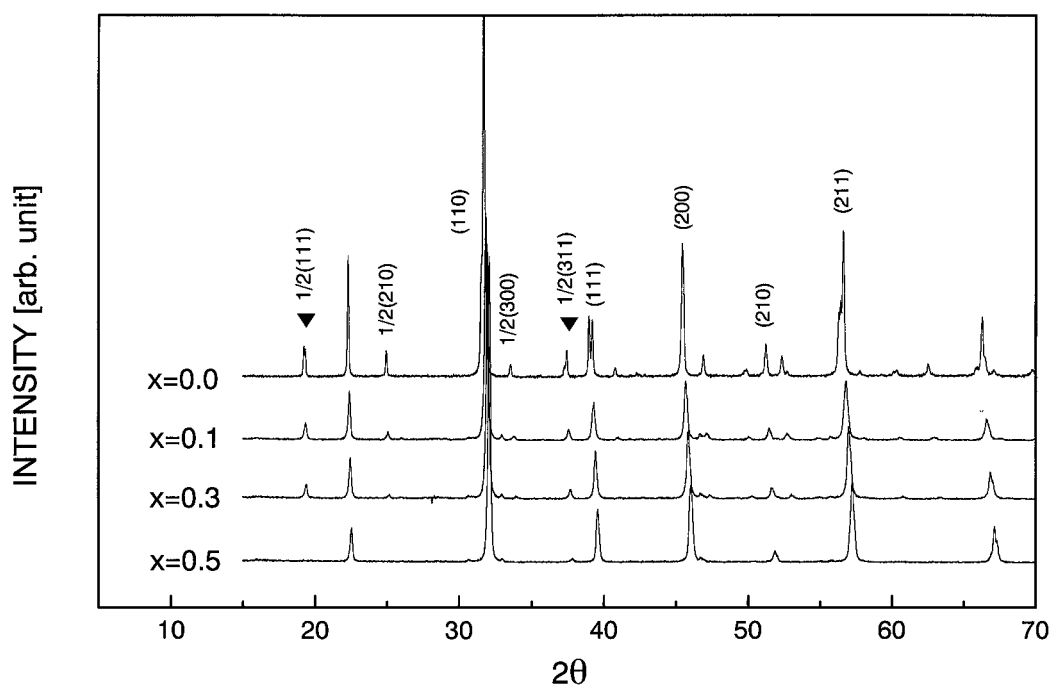


Figure 3 X-ray diffraction pattern of  $(1-x)\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3 - x\text{SrTiO}_3$  ceramics sintered at  $1500^\circ\text{C}$  for 4 h.

peak existed up to  $x = 0.5$ . Since the 1:1 ordering disappeared as  $x$  exceeded 0.3, the existence of  $1/2(311)$  extra peak in the specimen with  $x > 0.3$  is explained by the anti-phase tilting of octahedra. This result implies that the  $1/2(311)$  extra peak in LMN is associated with both the 1:1 ordering and the anti-phase tilting of the octahedra. The extra peaks for the in-phase tilting of the octahedra were not found in XRD pattern. However, since the in-phase tilting and the anti-parallel shift of A-site cations occurred simultaneously [7], the in-phase tilting was believed to exist in LMN and TEM results which will be shown in later showed the presence of in-phase tilting.

Rietveld refinement was carried out to determine the crystal structure of LMN. The refined results are summarized in Table I. First, the rhombohedral  $R\bar{3}$  model suggested by Lin *et al.* was tested [12]. As shown in

TABLE I Summary of Rietveld refinement for  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$  ceramics

| Space group | $R_p$ | $R_{wp}$ | $R_e$ | $\chi$ |
|-------------|-------|----------|-------|--------|
| $R\bar{3}$  | 14.09 | 17.87    | 4.41  | 4.06   |
| $Pbnm$      | 9.07  | 12.69    | 4.39  | 2.89   |
| $P2_1/n$    | 7.34  | 10.06    | 4.40  | 2.28   |

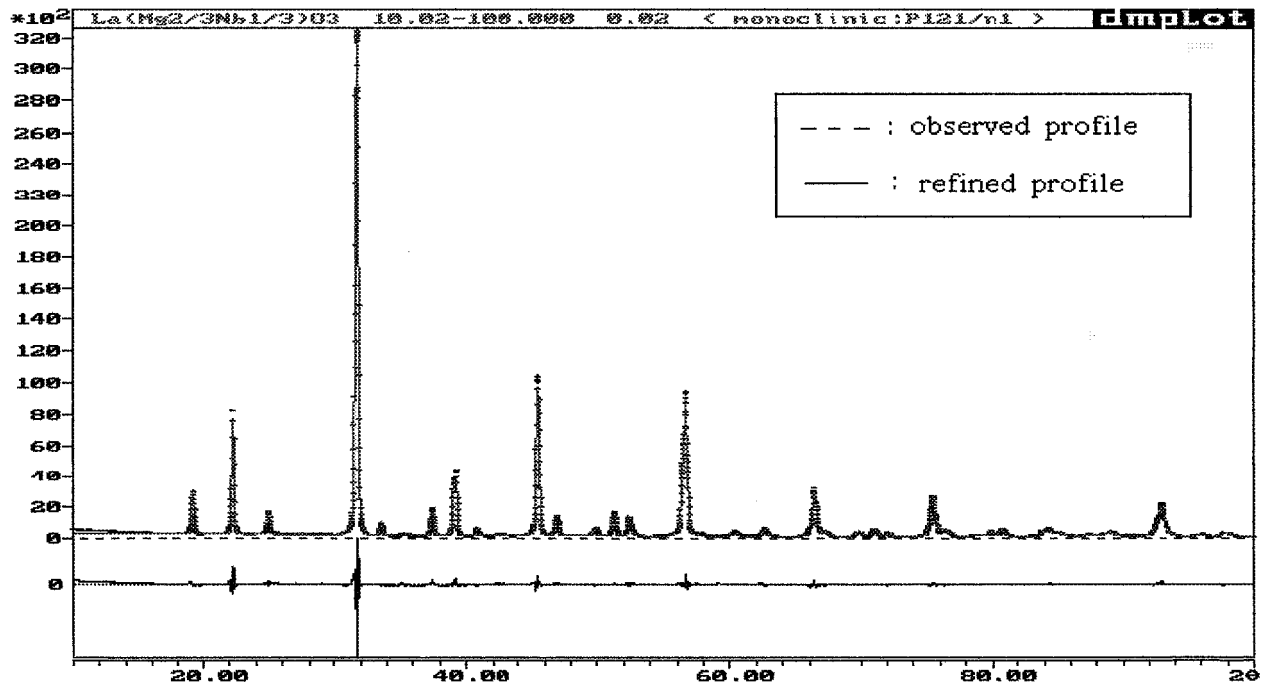
$\chi$  : Goodness of fit.

$$\chi^2 = (R_{wp}/R_{exp})^2$$

Table I.  $R$ -factors produced by this model were high. Moreover, in rhombohedral  $R\bar{3}$  structure, the in-phase tilting of octahedra which exists in LMN is not allowed [3, 4]. Therefore, LMN can't have the rhombohedral  $R\bar{3}$  structure. The orthorhombic structure  $Pbnm$  was also tested.  $R$ -factors of this structure were low as listed in Table I. However, the 1:1 chemical ordering is not

TABLE II Refined lattice parameters, fractional atomic coordinates, thermal parameters and occupancies of  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$  ceramics

| Lattice parameter  |           |           |           |           |           |
|--|-----------|-----------|-----------|-----------|-----------|
| $a = 5.6004(4) \text{ \AA}$ $b = 5.6414(4) \text{ \AA}$ $c = 7.9346(5) \text{ \AA}$ $\beta = 89.9819(3)^\circ$ |           |           |           |           |           |
| Positional and thermal displace parameters from X-ray refinement   |           |           |           |           |           |
| Atoms  | $x$       | $y$       | $z$       | $B$       | Occup.    |
| La   | 0.5046(3) | 0.5331(1) | 0.2511(2) | 2.1841(6) | 1.00000   |
| Mg(B')   | 0.00000   | 0.50000   | 0.00000   | 1.1015(2) | 0.50000   |
| Mg(B'')  | 0.50000   | 0.00000   | 0.00000   | 1.4384(4) | 0.0983(4) |
| Nb(B'')  | 0.50000   | 0.00000   | 0.00000   | 1.4384(4) | 0.4017(4) |
| O <sub>1</sub>   | 0.2901(6) | 0.2725(6) | 0.0326(2) | 0.9239(9) | 1.00000   |
| O <sub>2</sub>   | 0.2427(4) | 0.7814(3) | 0.0707(1) | 0.9239(9) | 1.00000   |
| O <sub>3</sub>   | 0.4362(8) | 0.9821(6) | 0.2543(5) | 0.9239(9) | 1.00000   |


 Figure 4 Observed and refined X-ray diffraction pattern of  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$ .

allowed in this structure. Therefore, the extra peak at  $1/2(111)$  position (odd-odd-odd reflection) should be explained by the anti-phase tilting of octahedra [3, 4]. However, according to the work of Glazer, the lowest-order reflection due to the anti-phase tilting is  $1/2(311)$  [3, 4]. Moreover, the behavior of the extra peak at  $1/2(111)$  position shown in Fig. 2 is different from that of  $1/2(311)$  peak. Therefore, the extra peak at  $1/2(111)$  position can't be explained by the anti-phase tilting of octahedra and only the 1:1 chemical ordering can explain the existence of the  $1/2(111)$  extra peak. As a result, the orthorhombic  $\text{Pbmn}$  is not the structure of LMN.

Finally, monoclinic structure  $\text{P}2_1/n$  which has both the 1:1 chemical ordering and  $a^- a^- c^+$  (or  $a^- a^- a^+$ ) tilting axis was tested [18]. As shown in Table I,  $\text{P}2_1/n$  model produced the most successful refinement results. The observed and calculated X-ray diffraction pattern is shown in Fig. 4. Refined atomic parameters and interatomic distances are given in Tables II and III, respectively. The refined lattice parameters are  $a = 5.6004 \text{ \AA}$ ,  $b = 5.6414 \text{ \AA}$ ,  $c = 7.9346 \text{ \AA}$ , and  $\beta = 89.9819^\circ$ . The pseudocubic unit cell lengths of LMN which were used by Glazer to denote the tilting system are  $a_p = b_p = 0.3927 \text{ nm}$  and  $c_p = 0.3921 \text{ nm}$

 TABLE III Selected atomic distances ( $\text{\AA}$ ) of  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$  ceramics

| MgO <sub>6</sub> -octahedron |        |    | LaO <sub>8</sub> -polyhedron |        |    |
|------------------------------|--------|----|------------------------------|--------|----|
| Mg-O <sub>1</sub>            | 2.0868 | 2× | La-O <sub>1</sub>            | 2.5711 | 1× |
| Mg-O <sub>2</sub>            | 2.1639 | 2× |                              | 2.7372 | 1× |
| Mg-O <sub>3</sub>            | 1.9849 | 2× |                              | 2.7547 | 1× |
| Mean                         | 2.0785 |    | La-O <sub>2</sub>            | 2.4359 | 1× |
| Shannon                      | 2.1    |    |                              | 2.4824 | 1× |
|                              |        |    |                              | 3.0508 | 1× |
| NbO <sub>6</sub> -octahedron |        |    | La-O <sub>3</sub>            |        |    |
| Nb-O <sub>1</sub>            | 1.9522 | 2× |                              | 2.4853 | 1× |
| Nb-O <sub>2</sub>            | 1.9777 | 2× | Mean                         | 2.5620 | 1× |
| Nb-O <sub>3</sub>            | 2.0514 | 2× | Shannon                      | 2.6350 |    |
| Mean                         | 1.9938 |    |                              | 2.54   |    |
| Shannon                      | 2.02   |    |                              |        |    |

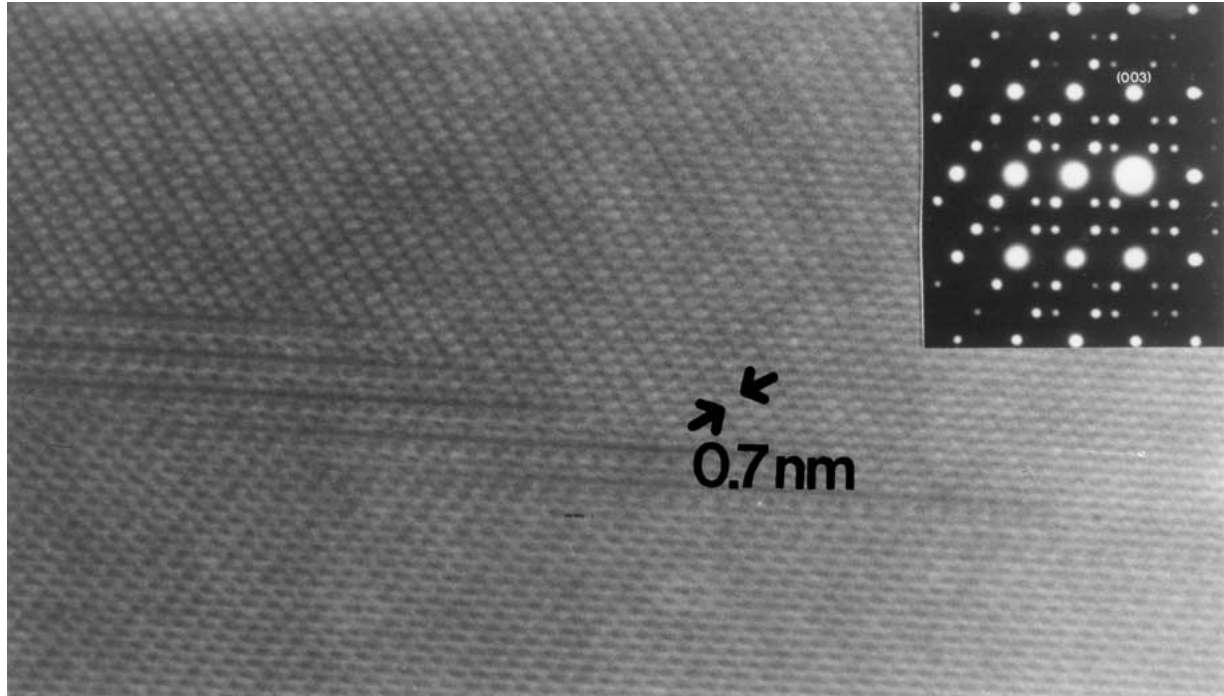
[3, 4]. Therefore, LMN has  $a^- a^- c^+$  tilting system. The X-ray diffraction pattern of LMN is the same as those of BLMN with  $x > 0.7$ . Thus, it can be concluded that BLMN specimens with  $x > 0.7$  also have the 1:1 ordered monoclinic structure.

Transmission electron microscopy studies were performed to investigate in detail variations in the crystal structure of BLMN with La content. Fig. 5a to d

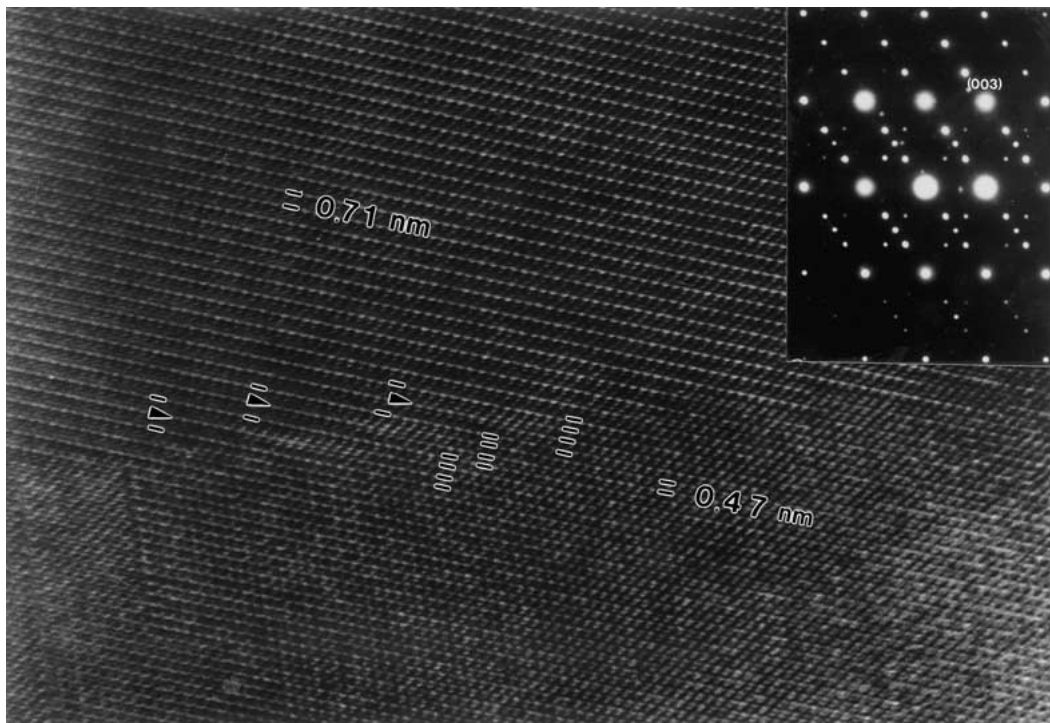
represent the HRTEM images and electron diffraction patterns of BLMN with  $x = 0, 0.1, 0.5$  and  $0.9$ , respectively. As seen Fig. 5a, the 1:2 ordered hexagonal structure was well developed in BMN, and that structure was found in the entire area of BMN. In the case of BLMN with  $x = 0.1$ , both the 1:2 and 1:1 ordered structures co-existed, as shown in Fig. 5b. Most of the area, however, has the 1:1 ordered cubic structure. The results shown in Fig. 5a and b indicate that a small amount of La substitution destroyed the 1:2 ordered hexagonal structure

and caused a 1:1 ordered cubic structure. Fig. 5c shows the HRTEM image and diffraction pattern of the specimen with  $x = 0.5$ . Only the 1:1 ordered cubic structure was found in that specimen. Similar results were obtained for the specimens with  $0.1 < x \leq 0.7$ .

Fig. 5d shows the HRTEM image of the [100] zone axis obtained from the BLMN with  $x = 0.9$ . The inset represents the [100] electron diffraction pattern taken from the same area. The extra reflections at  $1/2\{210\}$  and  $1/2\{300\}$  which are due to the anti-parallel shift of

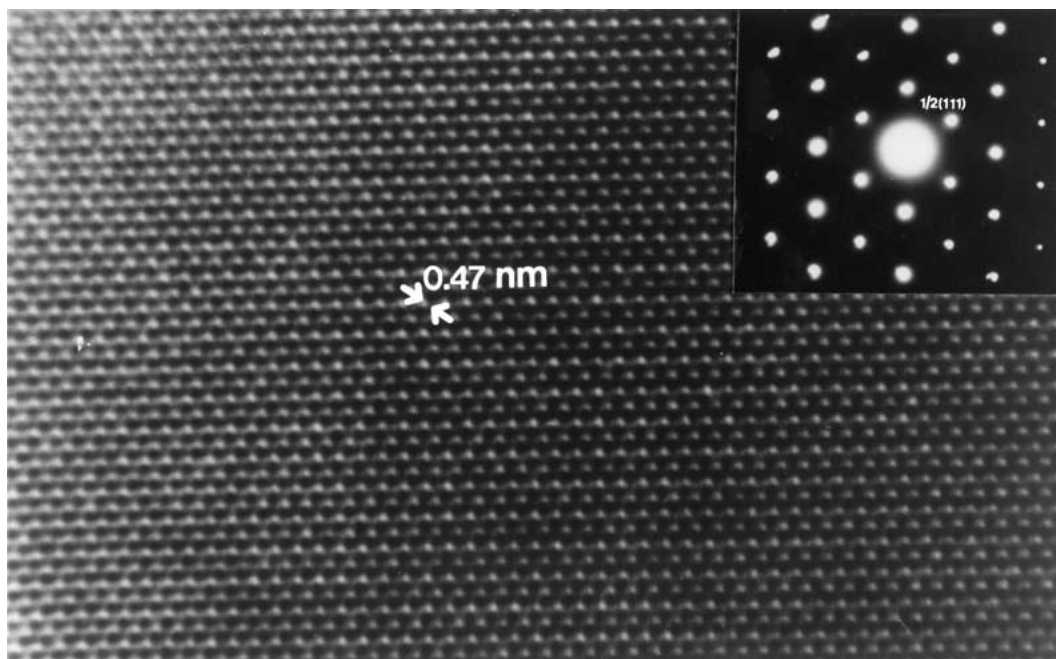


(a)

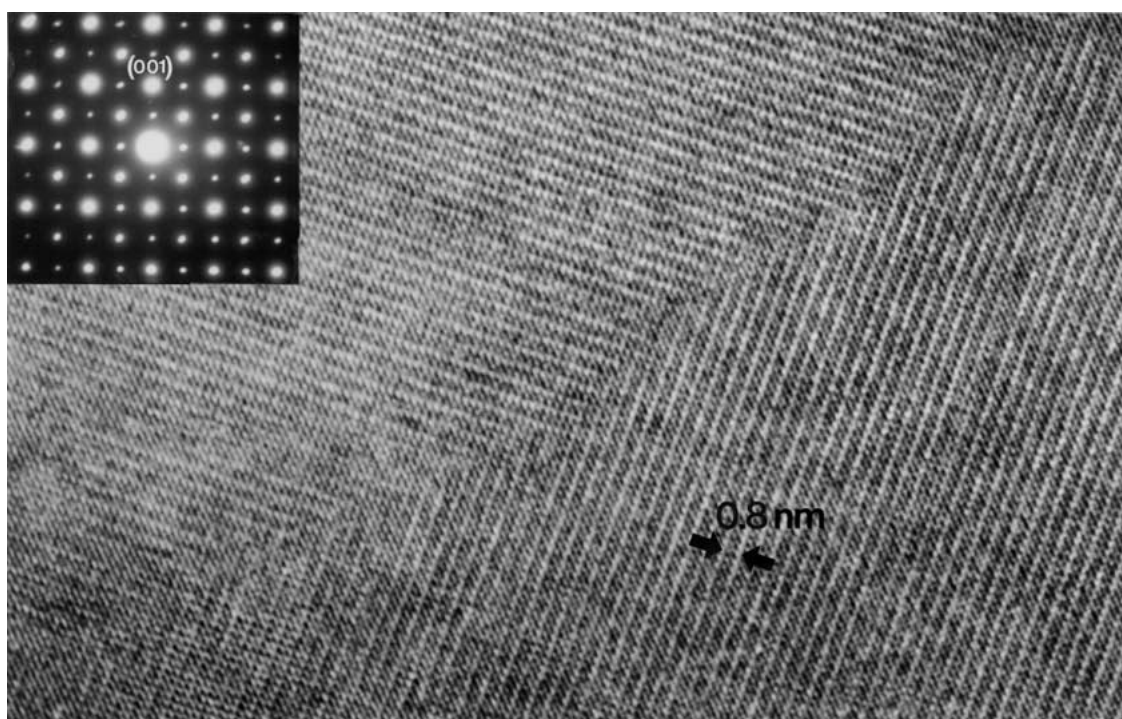


(b)

Figure 5 HRTEM images and electron diffraction patterns of (a)  $\text{Ba}(\text{Mg}_{1/3}\text{Nb}_{2/3})\text{O}_3$ , (b)  $(\text{Ba}_{0.9}\text{La}_{0.1})(\text{Mg}_{1.1/3}\text{Nb}_{1.9/3})\text{O}_3$ , (c)  $(\text{Ba}_{0.5}\text{La}_{0.5})(\text{Mg}_{1/2}\text{Nb}_{1/2})\text{O}_3$  and (d)  $(\text{Ba}_{0.1}\text{La}_{0.9})(\text{Mg}_{1.9/3}\text{Nb}_{1.1/3})\text{O}_3$ . (Continued.)



(c)

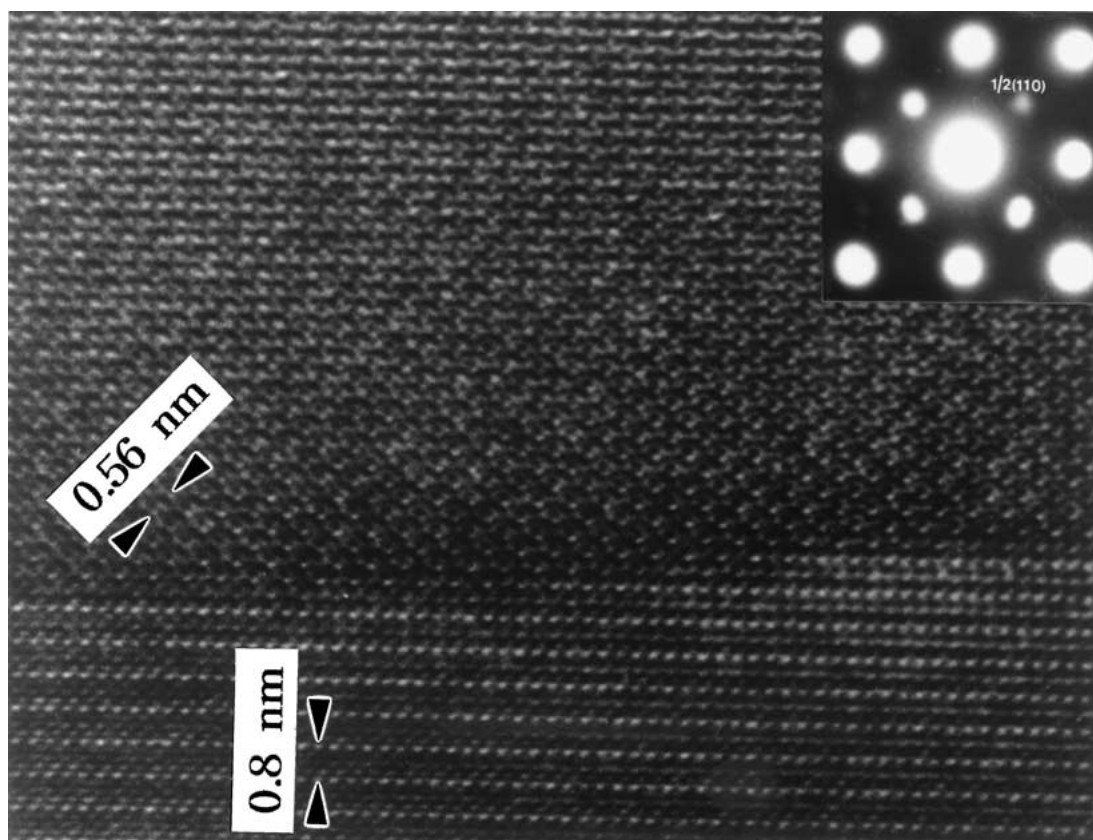


(d)

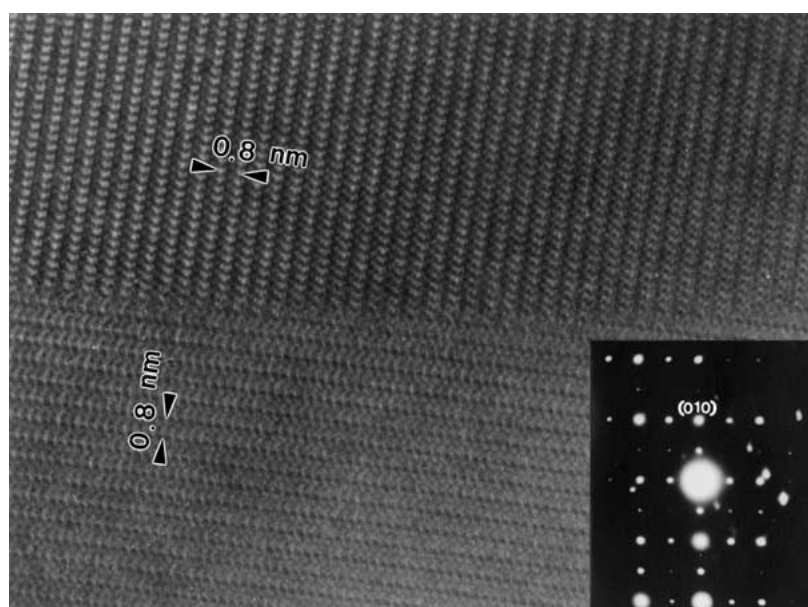
Figure 5 (Continued).

the A-site cations were observed in this pattern. Moreover, the  $1/2\{100\}$   $1/2\{100\}$  superlattice reflections that were not observed in XRD pattern were obtained in electron diffraction pattern. The HRTEM image also shows the new type of modulation along the  $[010]$  and  $[001]$  directions with the wavelength of 0.8 nm. According to the previous work, the  $1/2\{100\}$  and  $1/2\{110\}$  superlattice reflections were associated with the anti-parallel shift of the A-site cations and the in-phase tilting of the oxygen octahedra, respectively. Therefore, TEM analysis clearly showed the presence of the in phase tilting of oxygen octahedra in BLMN with  $x > 0.7$ .

Fig. 6a shows the HRTEM image with  $[001]$  zone axis taken from LMN specimen. The inset represents the electron diffraction pattern taken from the same area. The diffraction pattern shows the  $1/2\{110\}$  superlattice reflections indicating the presence of the in phase tilting of oxygen octahedra in LMN. Those superlattice reflections give rise to modulation along the  $[110]$  directions, with a wavelength of 0.56 nm in the HRTEM image shown in Fig. 6a. Fig. 6b shows the  $[100]$  zone axis HRTEM image and electron diffraction pattern. The diffraction pattern exhibits  $1/2\{100\}$  superlattice reflections in two directions. Those  $1/2\{100\}$  superlattice reflections give rise to two independent set of



(a)



(b)

Figure 6 HRTEM image and electron diffraction pattern of LMN with (a) [001] zone axis and (b) [100] zone axis.

$1/2\{100\}$  lattice fringe, seen in HRTEM image. The  $1/2\{100\}$  superlattice reflections in different directions therefore arose from the different variants.

#### 4. Conclusions

Variation in the crystal structure of  $(\text{Ba}_{1-x}\text{La}_x)(\text{Mg}_{(1+x)/3}\text{Nb}_{(2-x)/3})\text{O}_3$  with  $0 \leq x \leq 1$  was investigated in the present study. A 1:2 ordered hexagonal structure was found in the whole region of BMN. With increased La content, the structure changes to the 1:1 ordered cubic. When  $x = 0.1$ , both 1:1 and 1:2 or-

dered structures were found but most of the area had the 1:1 ordered cubic structure. The 1:1 ordered cubic structure was maintained up to  $x = 0.7$ . When  $x$  exceeded 0.7, however, BLMN is transformed from a 1:1 ordered cubic to a 1:1 ordered monoclinic structure. The 1:1 ordered monoclinic BLMN had both the in-phase and anti-phase tilting of oxygen octahedra.  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$  has the 1:1 ordered monoclinic  $P2_1/n$  structure. The lattice parameters of  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$  are  $a = 5.6004 \text{ \AA}$ ,  $b = 5.6414 \text{ \AA}$ ,  $c = 7.9346 \text{ \AA}$ , and  $\beta = 89.9819^\circ$ . The 1:1 ordered monoclinic  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$  is distorted by the in-phase

and the anti-phase tilting of octahedra with the  $a^- a^- c^+$  tilting system. The anti-parallel shift of A-site cations was also found in  $\text{La}(\text{Mg}_{2/3}\text{Nb}_{1/3})\text{O}_3$ .

## Acknowledgments

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